

Batch Adsorption: Intraparticle Adsorbate Concentration Profile Models

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Several models proposed describe the adsorbate concentration profile in an adsorbent particle. A theoretical consideration is made in this work to see what bulk concentration-time curves are predicted by these models when applied to a batch adsorption. A similar, but experimental, comparison was made in terms of parameter estimation by Kaguei et al. (1989).

Mass balance for batch adsorption with a linear isotherm is expressed:

$$\frac{\partial A}{\partial \tau} = \frac{1}{x^2} \frac{\partial}{\partial x} \left(x^2 \frac{\partial A}{\partial x} \right) \quad (1a)$$

$$\frac{dA_b}{d\tau} = -N \left(\frac{\partial A}{\partial x} \right)_{x=1} \quad (1b)$$

$$\text{at } \tau = 0, A = 0, A_b = 1 \quad (1c)$$

$$\text{at } x = 1, \frac{\partial A}{\partial x} = Bi(A_b - A) \quad (1d)$$

and λ_n is an n th positive root of the following equation:

$$\lambda \cot \lambda - 1 = \frac{\lambda^2}{N - \frac{\lambda^2}{Bi}} \quad (2d)$$

N th-Order Concentration Profile Model

Do and Mayfield (1987) assumed

$$A = a_0 + a_n x^n \quad (3a)$$

and found the exponent to be approximately:

$$n = 0.123\tau^{-0.68} \quad (3b)$$

Their study on the batch adsorption with an infinitely large volume of tank is modified to that with a finite volume of tank. The loss of adsorbate from the bulk fluid equals the gain of the adsorbent:

$$\frac{dA_b}{d\tau} = -\frac{N}{3} \frac{d\bar{A}}{d\tau} \quad (4a)$$

where \bar{A} is the average concentration in a particle.

$$\bar{A} = 3 \int_0^1 Ax^2 dx \quad (4b)$$

An integration then gives:

$$A_b - 1 = -\frac{N\bar{A}}{3} \quad (4c)$$

$$= -N \left(\frac{a_0}{3} + \frac{a_n}{n+3} \right) \quad (4d)$$

Exact Solution (Do and Rice, 1986)

$$(A_b)_{\text{exact}} = (A_b)_{\infty} + \sum_{n=1}^{\infty} A_n \exp(-\lambda_n^2 \tau) \quad (2a)$$

where

$$(A_b)_{\infty} = \frac{3}{3+N} \quad (2b)$$

$$A_n = \left[1 + \frac{N}{2} \left(1 - \frac{\lambda_n^2}{NBi} \right)^2 \left(\frac{1}{\sin^2 \lambda_n} - \frac{\cot \lambda_n}{\lambda_n} \right) \right]^{-1} \quad (2c)$$

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Combining Eq. 4d with Eqs. 1b and 1d in which Eq. 3a is substituted, one obtains:

$$\left(\frac{1}{Bi} + \frac{1}{n+3}\right) \frac{dA_b}{d\tau} = 3 - (3+N)A_b \quad (5)$$

Hence, let us approximate the power “-0.68” by “-2/3” so that Eq. 3b becomes (in the range of $\tau = 0.001$ -0.2):

$$n = 0.133\tau^{-2/3} \quad (6)$$

As a matter of fact, with the aid of Eq. 6, Eq. 5 is analytically solved to give:

$$(A_b)_{nth} = (A_b)_\infty \left\{ 1 + \frac{N}{3} \exp \left[- \left(1 + \frac{N}{3} \right) \psi \right] \right\} \quad (7a)$$

where

$$\psi = \frac{9}{1 + \frac{3}{Bi}} \left\{ \tau + \gamma^2 Bi \left[\tau^{1/3} - \gamma \tan^{-1} \left(\frac{\tau^{1/3}}{\gamma} \right) \right] \right\} \quad (7b)$$

$$\gamma = \sqrt{\frac{0.133}{3 + Bi}} \quad (7c)$$

Some other models examined in our note are discussed below.

Parabolic Concentration Profile Model

Liaw et al. (1979) and Rice (1982) assumed

$$A = a_0 + a_2 x^2 \quad (8a)$$

and obtained

$$(A_b)_{par} = (A_b)_\infty \left[1 + \frac{N}{3} \exp \left(- \frac{3+N}{\frac{1}{Bi} + \frac{1}{5}} \tau \right) \right] \quad (8b)$$

Quartic Concentration Profile Model

Do and Rice (1986) assumed

$$A = a_0 + a_2 x^2 + a_4 x^4 \quad (9a)$$

and obtained

$$(A_b)_{quar} = (A_b)_\infty + D_1 \exp(-d_1 \tau) + D_2 \exp(-d_2 \tau) \quad (9b)$$

where

$$d_1, d_2 = \frac{u_1}{2u_2} \left(1 \pm \sqrt{1 - \frac{4u_2}{u_1^2 (A_b)_\infty}} \right) \quad (9c)$$

$$D_i = \frac{1 - (A_b)_\infty [1 + 3Nd_i/140]}{2 - (A_b)_\infty u_1 d_i} \quad (i = 1, 2) \quad (9d)$$

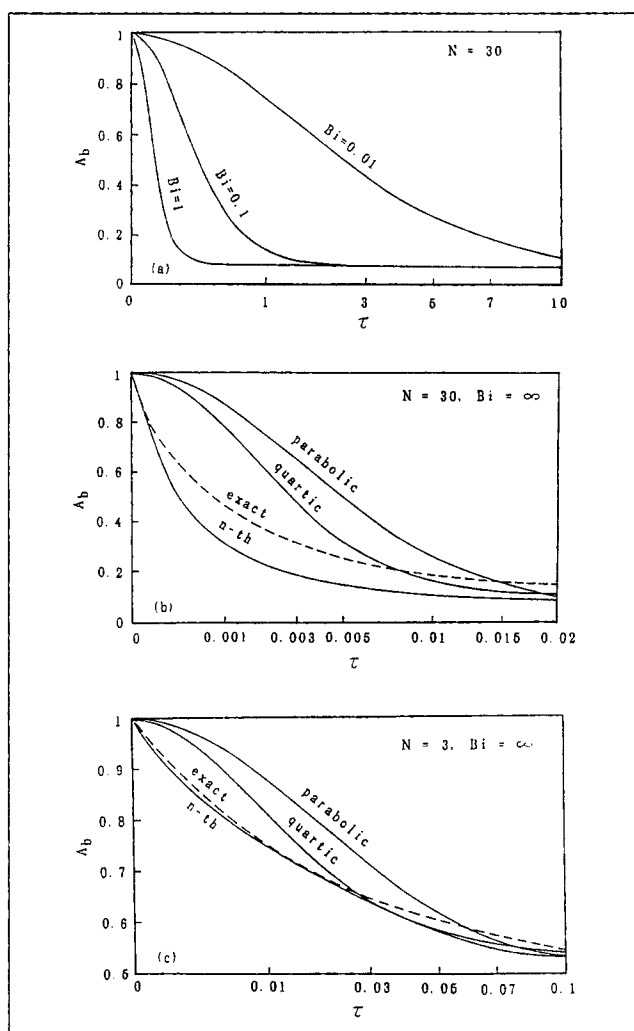


Figure 1. Exact solution $(A_b)_{exact} - \tau$ vs. $A_b - \tau$ curves for n th-order, parabolic and quartic concentration models: (a) $N = 30$, $Bi \leq 1$; (b) $N = 30$, $Bi = \infty$; (c) $N = 3$, $Bi = \infty$.

$$u_1 = \frac{11}{84} + \frac{3}{140} N + \frac{1}{3Bi} \quad (9e)$$

$$u_2 = \frac{1}{420} + \frac{9}{420Bi} \quad (9f)$$

Figures 1a-1c (abscissa proportional to $\sqrt{\tau}$) illustrate $A_b - \tau$ curves calculated from Eqs. 2a, 7a, 8b, and 9b, respectively. As shown in Figure 1a, when $Bi \leq 1$, there is no appreciable difference between all the curves, $(A_b)_{par} - \tau$, $(A_b)_{quar} - \tau$, $(A_b)_{nth} - \tau$ and $(A_b)_{exact} - \tau$. Figures 1b and 1c show the curves with $Bi = \infty$: When $N = 30$, all the curves for concentration profile models differ considerably from the $(A_b)_{exact} - \tau$ curve. The deviation becomes larger increasingly in the order of that with $(A_b)_{nth} - \tau$, $(A_b)_{quar} - \tau$, and $(A_b)_{par} - \tau$. When $N = 3$, the $(A_b)_{nth} - \tau$ curve is close to the $(A_b)_{exact} - \tau$ curve, but the other two curves differ from the $(A_b)_{exact} - \tau$ curve. The deviation is again the largest with $(A_b)_{par} - \tau$ curve.

It is obvious that the intraparticle concentration profile becomes less significant when Bi is small: fluid-to-particle mass-

transfer controls the overall rate. On the other hand, the concentration profile plays an important role when Bi is large: the overall rate is controlled by the diffusion of adsorbate in particle. In the latter case, as far as the prediction of bulk concentration in a batch tank is concerned, none of the concentration profile models so far proposed are satisfactory.

Rice and Do will propose a new approximation for the situation when $Bi \rightarrow \infty$ in their book titled *Applied Mathematics and Modeling for Chemical Engineers*, which will be published in 1994.

Notation

- A = dimensionless intraparticle concentration, $q/(KC_0)$
- A_b = dimensionless bulk fluid concentration, C/C_0
- Bi = Biot number, $k_m R/(KD_s)$
- C = bulk fluid concentration
- C_0 = initial bulk fluid concentration
- D_s = surface diffusivity
- K = linear equilibrium adsorption constant
- k_m = particle-fluid mass-transfer coefficient
- m_p = mass of adsorbent
- N = $3m_p K/(V\rho_p)$
- q = adsorbed concentration
- R = radius

- t = time
- V = liquid volume
- x = dimensionless radial distance variable

Greek letters

- ρ_p = particle density
- $\tau = D_s t/R^2$

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